

# Epitaxially Intergrown Conformational Polymorphs and a Mixed Water/Methanol Solvate of 5′-Deoxy-5′-iodoguanosine

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**S** [Supporting Information](#page-9-0)

ABSTRACT: 5′-Deoxy-5′-iodoguanosine (I) crystals deposited from mixtures of water and methanol grow as nonsolvated hybrids of conformational polymorphs (Ia, Ib) and as a mixed solvate (Ic). Some solvent-free crystals are purely Ia, while others have varying amounts of Ib epitaxially intergrown with Ia. In Ia and Ib the conformations differ primarily by torsion about the C4′−C5′ bond (guanosine numbering scheme), which dramatically affects the iodine atom position. Powder diffraction and reconstructed reciprocal-lattice-slice images had small peaks incompatible with Ia. Some solvent-free crystals



required lattices for both Ia and Ib to index all observable reflections. Unit-cell dimensions for Ia and Ib suggest the potential for epitaxial intergrowth. Hydrogen-bond networks in Ia and Ib are essentially identical and result in double layers of molecules in the ab plane, with layers of iodine at the layer surfaces. The iodine layers of Ia and Ib are incompatible: in Ia adjacent iodine atom layers interdigitate slightly, whereas in Ib they do not. Theoretical calculations support the conclusion that at room temperature Ia is the thermodynamically more stable polymorph and that Ib represents a kinetic product.

# 1. INTRODUCTION

Polymorphism is the well-known phenomenon in which a compound crystallizes with more than one distinct crystal structure but identical chemical composition.<sup>[1](#page-9-0)−[3](#page-9-0)</sup> Broadly speaking, there are two main types: polymorphism that results from packing differences only, and polymorphism arising from differences in molecular conformation. Crystals containing different solvents or quantities of solvent are not polymorphs but have been referred to as "solvatomorphs" and "pseudopolymorphs", though not without controversy.[4](#page-9-0)−[6](#page-9-0) Polymorphs and related solvates can have commercial ramifications in crystal engineering because changes in crystal structure, whether due to packing differences, conformational changes, or both, can lead to dramatic differences in physical properties. In the pharmaceutical industry, for example, polymorphs in drug formulations may exhibit very different dissolution properties, which can affect bioavailability. The most famous example is probably the antiretroviral drug ritonavir.<sup>[7](#page-9-0)</sup> In the field of organic electronics, properties of crystals and crystalline thin films are critically dependent upon the manner in which electronically active functional groups on adjacent molecules interact. Even subtle changes in crystal packing can dramatically affect properties such as conductivity and electron/hole mobility.<sup>[8](#page-9-0),</sup>

5′-Deoxy-5′-iodoguanosine (I, Scheme 1) is a versatile starting material used to produce a wide variety of 5′-





 $deoxyguanosines<sup>10</sup>$  $deoxyguanosines<sup>10</sup>$  $deoxyguanosines<sup>10</sup>$  for use in site-specific modification of  $RNA<sup>[11-13]</sup>$  $RNA<sup>[11-13]</sup>$  $RNA<sup>[11-13]</sup>$  $RNA<sup>[11-13]</sup>$  $RNA<sup>[11-13]</sup>$  The synthesis of I has been reported via several synthetic routes.<sup>[14](#page-9-0)–[16](#page-9-0)</sup> Its crystal structure, however, has thus far not been reported, perhaps as a consequence of the unusual diffraction encountered in the present study. In the course of work on substituted 5′-deoxyguanosines, we first obtained very small rectangular plate-like crystals of I by evaporation of MeOH from a solution in mixed  $H<sub>2</sub>O/MeOH$ . These crystals contained two independent molecules of I that had quite different conformations, along with water and methanol. Subsequent crystals grown from water and or water/methanol

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 ${}^a$ Formula does not include the diffuse feature that was modeled as partial occupancy disordered water/methanol.  ${}^b$ Includes contribution from the diffuse feature that was modeled as partial occupancy disordered water/methanol.

were mostly lath-shaped and solvent free. Their structure, Ia, solved easily, but the diffraction pattern contained unusual features. Comparison of powder diffraction patterns for the solvent-free crystals to a powder pattern calculated from the refined single-crystal structure revealed a few small peaks in the experimental pattern that indicated the presence of a phase that did not fit either Ia or that of the solvate. Similarly, reciprocallattice-slice images calculated from a full set of raw diffraction frames revealed a small number of weak Bragg peaks that do not correspond to the structure of Ia. The unusual diffraction from these crystals prompted a closer look. At the Advanced Light Source (ALS, LBNL, Berkeley, CA), a miniscule squat slab-shaped crystal gave diffraction that clearly showed two different but related mutually oriented reciprocal lattices. One unit cell matched that of Ia, while the other was for a second polymorph, Ib.

Epitaxial intergrowth of polymorphs and solvates is not unknown in molecular crystals, but it is not common. Two examples of polymorphic hybrid molecular crystals are DL-homocysteine thiolactone hydrochloride<sup>[17](#page-9-0)</sup> and aspirin,<sup>[18](#page-9-0)</sup> but these are strictly packing polymorphs, whereas Ia and Ib are conformational polymorphs. An example of epitaxial intergrowth of anhydrous and hydrated forms with biological implications is uric acid.<sup>[19](#page-9-0),[20](#page-9-0)</sup> In this paper we describe the structures and hydrogen bonding in polymorphs Ia and Ib, and suggest plausible reasons for the observed variable intergrowth of the two forms. For the sake of completeness, we also report the crystal structure of the mixed  $H_2O/MeOH$  solvate Ic.

# 2. EXPERIMENTAL SECTION

2.1. Synthesis and Crystallization. Synthesis of I followed the procedure of McGee and Martin.<sup>[14](#page-9-0)</sup> The first crystals were produced by dissolving I (50 mg) in boiling methanol/water (3 mL, 1:1) and allowing the methanol to evaporate slowly at RT. The initial preparation gave crystals of the  $H<sub>2</sub>O/MeOH$  solvate Ic. IR (Nujol): 3386, 3283, 3136, 1704, 1678, 1628, 1565, 1529, 1349, 1330, 1244, 1161, 1120, 1008, 941, 826, 683. Deuterated crystals were prepared by dissolving I (50 mg) in boiling  $D_2O$  (5 mL) and continuing the heating in a boiling water bath for 5 min. The heat source was turned off, and the solution was slowly cooled to RT. The crystals were filtered, washed with  $D_2O$ , and then either air-dried or kept wet. The resulting crystals did not contain solvent and was indexed to give cell dimensions of form Ia. Subsequent preparations from either methanol/water or water alone (deuterated or protonated) appeared to give only the nonsolvated crystals (i.e., Ia and hybrid Ia/Ib), though not all batches were subject to crystallographic analysis. IR (Nujol): 2597, 2507, 2444, 2351, 1665, 1650, 1563, 1537, 1456, 1346, 1293, 1261,1082, 1070, 1015, 958, 923, 779, 706. The <sup>1</sup> H-NMR spectrum in wet deuterated DMSO matched the data of McGee and Martin.<sup>[14](#page-9-0)</sup>

2.2. Single-Crystal Data Collection. All crystals obtained in this work were very small. The nonsolvated crystals grew as flattened rods or laths with typical dimensions  $0.15 \times 0.04 \times 0.01$  mm, elongated along b and flattened along c, though a few were slightly longer and/or thicker (Table 1, [SI Figure S1](http://pubs.acs.org/doi/suppl/10.1021/acs.cgd.6b00981/suppl_file/cg6b00981_si_001.pdf)). A very few crystals were semiregular squat slabs, typically  $0.04 \times 0.03 \times 0.01$  mm. Crystals of the H<sub>2</sub>O/ MeOH solvate were rectangular plates with maximum dimension ∼0.20 mm, but only ∼0.005−0.010 mm thick. For analysis, crystals were submerged in a drop of polyisobutene oil, cleaned of adhering smaller crystals, and mounted on either fine glass fibers or on polyimide scoops (MiTeGen) under a polarizing microscope. Crystals of Ia and Ic were mounted directly into the 90.0(2) K cold-gas stream (CryoIndustries LT-2), and data were collected using  $\phi$  and  $\omega$  scans with Cu Kα X-rays on a Bruker-Nonius X8 Proteum diffractometer equipped with multilayer focusing optics. Diffraction data for Ib from a slab-shaped crystal containing both Ia and Ib were collected using  $\phi$ and  $\omega$  scans at 100(2) K (Oxford Cryosystems Cryostream 800 Plus) on beamline 11.3.1 at the ALS using a wavelength of 1.0333(1) Å on a Bruker Photon100 CMOS equipped D8 diffractometer. In all cases, raw diffraction images were integrated using Saint-Plus in APEX2.<sup>[21](#page-9-0)</sup> Tools within APEX2 for handling twinned crystal data were used to partition reflections from Ia and Ib for separate indexing. Data scaling, merging, and absorption correction were performed using SADABS.<sup>2</sup> For Ib, a postrefinement correction using  $XABS2<sup>23</sup>$  $XABS2<sup>23</sup>$  $XABS2<sup>23</sup>$  was beneficial, presumably because this alleviated some residual effects of reflection overlap for the superimposed diffraction maxima of Ia with those of Ib.

2.3. Structure Solution and Refinement. Structures were solved using SHELXT<sup>24</sup> and refined using SHELXL.<sup>[25](#page-9-0)</sup> In Ia, all hydrogen atoms were found in difference maps. Hydrogens attached to carbon or oxygen were included using riding models, but those of the

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exocyclic amine were refined. In Ib, carbon-bound hydrogens were found in difference maps, but those bound to nitrogen or oxygen were placed at calculated positions: all were modeled as riding. In Ic, most hydrogens were initially found in difference maps, but some of the solvent hydrogens were placed at geometrically plausible hydrogen bonding positions and treated as riding. For the riding models, standard low-temperature distances were applied [0.95 Å  $(C_{sp2}H)$ , 0.99 Å ( $R_2CH_2$ ), 1.00 Å ( $R_3CH$ ), 0.88 Å ( $N_{sp2}H$ ), and 0.84 Å ( $OH$ )]. In all cases hydrogen atom  $U_{\text{iso}}$  values were tied to their parent atom  $U_{\text{eq}}$  values (1.5 $U_{\text{eq}}$  for OH and CH<sub>3</sub>, 1.2 $U_{\text{eq}}$  for all others). Crystal data, data collection, and experimental details for Ia, Ib, and Ic are summarized in [Table 1.](#page-1-0)

2.4. Powder Diffraction. Nonsolvated crystals of I were ground with a small agate mortar and pestle. A small pellet (∼0.2 mm diameter) of compressed powder was held in a fine polymer loop by a thin film of light polyisobutene oil. Two-dimensional (2-D) powder diffraction frames were collected on the same Bruker-Nonius X8 Proteum diffractometer used for single-crystal studies of Ia and Ic, but using the finest aperture beam-tunnel available. Data collection frames consisted of 360 $^{\circ}$  rotation about the  $\phi$  axis for 300 s at RT. Frames were radially integrated using DataSqueeze<sup>[26](#page-9-0)</sup> or PhaseID within  $APEX2<sup>21</sup>$  $APEX2<sup>21</sup>$  $APEX2<sup>21</sup>$  Calculated diffraction patterns based on the single-crystal structures of Ia and Ib were obtained using CrystalDiffract from the CrystalMaker suite. $27$  For ease of comparison, calculated powder peaks were broadened to approximate the width of the experimental powder diffraction peaks.

2.5. Computational Details. Conformational energies for I in 5° increments of torsion angle about the glycosidic bond between purine and deoxyribose moieties, and for the iodine relative to deoxyribose were calculated using the NWChem software package.<sup>[28](#page-9-0)</sup> The torsion angle in question was fixed and the geometry was allowed to relax for each angle using the B3LYP functional.<sup>[29,30](#page-9-0)</sup> Cohesive energies between planes of iodine atoms in adjacent layers were calculated for the layer interfaces present in Ia and Ib using NWChem, with a dispersion corrected functional (B3LYP-D3).<sup>[31](#page-9-0)</sup> Pairwise interaction energy corrected functional (B3LYP-D3).<sup>31</sup> Pairwise interaction energy decomposition analysis (PIEDA)<sup>[32](#page-9-0)</sup> was performed using the GAMESS-US software package,  $33$  with the MP2 method.  $34$  To perform the calculations efficiently, while still properly describing the high number of electrons around iodine, two basis sets were employed at the same time. The lighter atoms, C, H, N, and O used the 6-31G\* basis set, while iodine used the larger 6-311G\* basis set.

# 3. RESULTS

3.1. General Structural Features of 5′-Deoxy-5′ iodoguanosine (I). Bond parameters for individual crystallographic models of I are largely unremarkable. The guanosine skeleton consists of two relatively rigid groups: a flat purine ring system and a deoxyribose ring with envelope geometry. The overall conformation of the molecule is determined by rotation about the glycosidic bond (N9−C1′), which orients the ring systems relative to each other. A less dramatic conformational degree of freedom consists in rotation about C4′−C5′, which defines the torsion of the iodine substituent at the C5′ position. Variation in puckering<sup>[35](#page-10-0)</sup> of the deoxyribose ring is also possible. The molecule contains several potential strong hydrogen-bond donors and acceptors on both deoxyribose and purine moieties, so it is not surprising that major features of the crystal packing in both solvent-free and solvated forms are determined by hydrogen-bonding interactions.

3.2. Major Polymorph (Ia). The crystal used to refine the structure of polymorph Ia appeared to be solely Ia. The lathshaped crystal was monoclinic, space group  $P2_1$ , with cell dimensions  $a = 7.6285(4)$  Å,  $b = 4.9951(2)$  Å,  $c = 16.1724(7)$ Å,  $\beta = 91.843(2)$ °. An ellipsoid plot is shown in Figure 1a. Deviation from planarity for the purine ring system is negligible  $(0.0105 \text{ Å})$ . In the deoxyribose ring, the flap atoms  $(C2', C3',$ C4′) form a dihedral angle with the envelope atoms (C2′, C1′,



Figure 1. Ellipsoid plots (50% probability) of (a) polymorph Ia and (b) polymorph 1b. In Ia the iodine is gauche relative to O1′ whereas in Ib it is anti.

O1′, C4′) of 37.0(3)°. The torsion angle (C8−N9−C1′−O1′) between the purine and deoxyribose rings is  $28.5(7)$ °, while the torsions that define the iodine position  $[O1'–C4'–C5'–I1 =$ 61.2(5)° and C3′–C4′–C5′–I1 = 178.8(4)°] are gauche and anti respectively. Some conformational parameters are given in [Table 2](#page-3-0).

In the crystal, all hydrogen bonding is intermolecular. A series of  $R_2^2(10)$  motifs join the hydroxyl groups of the deoxyribose moiety to the purine of a translationally  $(1 + x, 1+)$  $y, z$ ) equivalent molecule into chains parallel to (110) by O– H…N and O−H…O<sub>C=O</sub> hydrogen bonds. The carbonyl oxygen is also involved as a bifurcated acceptor in an  $R_2^1(6)$ motif involving the ring NH and one hydrogen (H1N2) of the exocyclic  $NH<sub>2</sub>$  group of a molecule related by the  $2<sub>1</sub>$  screw axis parallel to  $b$ . This same exocyclic NH<sub>2</sub> donor hydrogen (i.e., H1N2) is bifurcated: it also interacts with O2' of a  $2<sub>1</sub>$  screwrelated molecule (via  $1 - x$ ,  $-0.5 + y$ ,  $1 - z$ ). Crystal packing for Ia is shown in [Figure 2](#page-3-0)a, and hydrogen bond parameters are listed in [Table 3](#page-3-0). There are no  $\pi-\pi$  overlaps between purines on adjacent molecules. The dihedral angle between purine rings related by the  $2<sub>1</sub>$  screw axis within one unit cell (i.e., via  $1 - x$ ,  $-0.5 + y$ ,  $1 - z$ ) is 78.88(8)°. The net result of these strong hydrogen-bonding interactions are double layers of molecules that extend in the ab plane. These double-layered planes have exposed iodine atoms at the double-layer surfaces. In Ia the C5′−I1 bond is oriented such that the iodine atom protrudes *outside* the unit cell box and is  $0.4258(3)$  Å from the *ab* plane, leading to a corrugated interface at which iodine atoms on adjacent double layers interdigitate ([Figure 3a](#page-4-0)).

<span id="page-3-0"></span>Table 2. Conformation-Defining Torsion Angles (deg) and Deoxyribose Ring Puckering Parameters for Polymorphs Ia and Ib, and for the Mixed Water/Methanol Solvate Ic

	Ia	Ib	$\text{Ic}_{\text{a}}$	$\text{Ic}_\text{b}$
$C4-N9-C1' - O1'$ (deg)	$-152.4(5)$	$-151(2)$	$-154.1(7)$	$-76.2(9)$
$C8-N9-C1' - O1'$ (deg)	28.5(7)	33(3)	38.0(11)	100.8(9)
$Q1' - C4' - C5' - I1'$ (deg)	61.2(5)	168.8(11)	174.9(5)	$-54.8(9)$
$C3' - C4' - C5' - I1'$ (deg)	178.8(4)	$-76.8(19)$	$-66.8(9)$	66.2(9)
Cremer-Pople $Q(A)$	0.384(6)	0.37(2)	0.405(9)	0.397(9)
Cremer-Pople $\Phi$ (deg)	93.1(8)	100(3)	83.6(11)	268.2(12)



Figure 2. Crystal packing in (a) Ia and (b) Ib viewed as a projection onto the ac plane. Hydrogen bonding patterns in the double layers (thin blue lines) are ostensibly the same, which highlights the compatibility of the two polymorphs for epitaxial intergrowth. The  $R_2^2(10)$  and  $R_2^1(6)$  ring motifs are highlighted by thick semitransparent red and blue lines, respectively. The only substantial difference is the orientation of the iodine atoms (purple) at the top (gauche) and bottom (anti) of each image.

3.3. Minor Polymorph Ib. The crystal used to refine the structure of Ib was an epitaxially intergrown hybrid of polymorphs Ia and Ib. The structure is monoclinic, space group  $P2_1$ , with cell dimensions  $a = 7.5536$  (15) Å,  $b =$ 4.9860(8) Å,  $c = 17.691(3)$  Å,  $\beta = 102.172(12)$ °. An ellipsoid plot of Ib is shown in [Figure 1](#page-2-0)b. Bond lengths and angles are similar to the equivalent bonds in polymorph Ia, but are less precise because of the necessary deconvolution of overlapping reflections during data integration.<sup>[21](#page-9-0)</sup> The geometry of the guanosine moiety is very similar to that in Ia [torsion angle  $(C8-N9-C1'-O1') = 33(3)°$ , but the  $(O1'-C4'-C5'-I1)$ and  $(C3'-C4'-C5'-11)$  torsions of  $168.8(11)^\circ$  and −76.8(19)° define anti and (suboptimal) gauche conformations

Table 3. Hydrogen Bonding in Polymorphs Ia and Ib

polymorph Ia	$d_{\text{D}H}$ (A)	$d_{\text{H}\dots A}$ (A)	$d_{\text{D}A}$ (A)	$\n  DHA\n$ $(\text{deg})$		
$N1-H1\cdots O1^{ia}$	0.88	2.03	2.845(6)	153.1		
$N2-H1N2\cdots O1^{ia}$	0.82(5)	2.28(6)	3.001(7)	147(8)		
$N2-H1N2\cdots O2'$	0.82(5)	2.37(8)	2.944(8)	127(8)		
$N2-H2N2\cdots N3^{iia}$	0.82(5)	2.46(6)	3.151(9)	143(7)		
$O2' - H2'1 \cdots O1$ <sup>iiia</sup>	0.84	2.04	2.811(6)	151.3		
$O3' - H3'1 \cdots N7$	0.84	1.98	2.774(6)	158.5		
Symmetry codes: ia = $(-x + 2, y - 1/2, -z + 1)$ , iia = $(-x + 1, y - 1/2, -z)$ $+$ 1), <i>iiia</i> = $(x - 1, y - 1, z)$						
polymorph (Ib)	$d_{\text{D} \cdots \text{H}}$ (Å)	$d_{H\cdots A}$ (A)	$d_{\text{D}A}$ (A)	$\angle$ DHA (deg)		
$N1-H1\cdots O1^{ib}$	0.88	2.04	2.83(2)	149		
$N2 - H1N2 \cdots Q1^{ib}$	0.88	2.34	3.14(2)	151.1		
$N2-H1N2\cdots O2^{\prime iib}$	0.88	2.22	2.98(2)	144.1		
$N2-H2N2\cdots N3^{iib}$	0.88	2.44	3.03(2)	125.1		
$O2' - H2'1 \cdots O1^{iiib}$	0.84	2.02	2.77(2)	148.5		
$O3' - H3'1 \cdots N7^{iiib}$	0.84	1.95	2.75(2)	158.2		
Symmetry codes: $ib = (-x, y + 1/2, -z + 1)$ , $iib = (-x + 1, y + 1/2, -z + 1)$ , $iiib = (x + 1, y + 1, z)$						

respectively, for the iodine substituent, i.e., opposite to Ia. Conformation defining parameters for Ib are given in Table 2.

Hydrogen bonding in Ib is virtually identical to that in Ia. The same graph-set motifs are present (Figure 2b) with just a few minor differences in geometric parameters, which are given in Table 3. The dihedral angle between purine rings related by the 2<sub>1</sub> screw axis within one unit cell [i.e., via  $(1 - x, 1/2 + y, 1)$ ]  $(-z)$ ] is 77.8(3)°, similar to Ia. The main differences arise from consequences of the conformations of the iodine substituents. In Ib, the iodine atom is  $0.5187(12)$  Å from the *ab* plane on the inside of the unit cell box. This leads to a much less corrugated interface with no opportunity for iodines on adjacent double layers to interdigitate ([Figure 3](#page-4-0)b).

3.4. Water/Methanol Solvate (Ic). Crystals of the water/ methanol solvate are orthorhombic, space group  $P2_12_12_1$ , with cell dimensions  $a = 9.2724(2)$  Å,  $b = 10.5929(3)$  Å,  $c =$ 32.2592(8) Å. The asymmetric unit for Ic contains two independent molecules,  $Ic_a$  and  $Ic_b$ , and is shown in [Figure 4](#page-4-0).

As with the solvent-free structures, there are no unusual bond lengths or angles. The conformations of I in  $Ic_a$  and  $Ic_b$ , however, are quite different from each other, and from Ia, though  $Ic<sub>a</sub>$  is similar to Ib. An overlay of the four conformations is shown in [Figure 5](#page-4-0). For  $Ic_{a}$ , the torsion between purine and deoxyribose moieties is  $38.0(11)^\circ$ , whereas in Ic<sub>b</sub> it is  $100.8(9)$ °. The torsion of the iodine substituent relative to the ring oxygen of deoxyribose in  $\text{Ic}_a$  is anti  $[174.9(5)^\circ]$ , whereas that of Ic<sub>b</sub> is gauche [−54.8(9)°], but in the opposite sense to that of Ib. Conformational parameters are summarized in Table 2. The incorporated solvent consists of an ordered methanol, three relatively well-ordered waters and an ill-defined elongated blob of electron density located in small regions

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Figure 3. Interlayer packing in (a) Ia and (b) Ib viewed along  $a$  and −a of the respective crystal structures. In Ia, layers of gauche-oriented iodine substituents interdigitate, whereas in Ib the anti-oriented layers of iodines abut each other. The relative I···I contact area is thus much greater in Ia than in Ib.



Figure 5. An overlay of the four experimentally determined conformations of I, based on a least-squares fit of atoms in the purine ring system. The conformations in polymorphs Ia and Ib differ primarily in the torsion of the iodine atom (far left). One conformation in the solvated crystals  $Ic<sub>a</sub>$  is very similar to that of polymorph Ib, while in the other,  $Ic_{b}$ , the deoxyribose ring position is substantially different.

between molecules  $Ic_a$  and  $Ic_b$  near the middle of the unit cell. We tried many different ways to model this feature, including various composites of disordered water, methanol, and mixtures of the two, none of which were truly satisfactory. Use of the SQUEEZE routine in Platon<sup>36</sup> to factor out this diffuse electron density was no better than a partial occupancy disordered solvent model. We settled with a composite of methanol and water (without hydrogen atoms), though its actual composition remains unknown. This has little bearing on the overall structure.

Hydrogen bonding within the solvate is considerably more complicated than Ia or Ib as it involves all the solvent molecules in addition to I. There are, however, several easily recognizable graph-set motifs: an  $R_2^2(8)$  ring joins inequivalent purines in the same asymmetric unit by pairs of N−H···O hydrogen bonds; an  $R_2^2(8)$  ring joins inequivalent purines in adjacent asymmetric units (via x, 1+ y, z) by pairs of long N− H···N hydrogen bonds; an  $R^3(8)$  ring involves the carbonyl (as a bifurcated acceptor) and exocyclic amine of an inequivalent purine moiety (same asymmetric unit), and a water molecule; an  $R_1^1(6)$  ring between inequivalent deoxyribose moieties related via the  $2<sub>1</sub>$  screw axis along b. The dihedral angle between purine moieties within one asymmetric unit is  $2.4(2)^{\circ}$ , so the propagation of both sets of  $R_2^2(8)$  motifs join the main molecules into gently undulating ribbons that extend parallel to b. Adjacent ribbons (via  $1 - x$ , 0.5 + y, 0.5 - z)  $\bar{x} - \pi$  stack



Figure 4. An ellipsoid plot (50% probability) of Ic. The conformations of the two molecules in the asymmetric unit Ic<sub>a</sub> on the left and Ic<sub>h</sub> on the right are quite different. A poorly defined feature that was modeled as a disordered superposition of partial occupancy water and methanol is omitted to improve clarity.

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<span id="page-5-0"></span>between the five- and six-membered rings of symmetry inequivalent pairs of purine moieties, with layer spacing 3.330(19) Å. These stacked bilayers of ribbons are built into an extended three-dimensional (3-D) structure by a complex network of O−H···N, N−H···O, and O−H···O hydrogen bonds (see [SI Table S1 and Figure S2\)](http://pubs.acs.org/doi/suppl/10.1021/acs.cgd.6b00981/suppl_file/cg6b00981_si_001.pdf) involving the main molecules, water, methanol, and the poorly defined disordered feature mentioned above.

3.5. Unusual Single-Crystal Diffraction from Solvent-Free I. Most crystals of solvent-free I were laths ([SI, Figure S1\)](http://pubs.acs.org/doi/suppl/10.1021/acs.cgd.6b00981/suppl_file/cg6b00981_si_001.pdf) that gave diffraction patterns that indexed as Ia. Some crystals gave diffraction patterns showing a single reciprocal lattice corresponding to Ia, but many others had faint additional reflections. These were consistently along  $c^*$  only, but had intensities that varied from crystal to crystal. Crystals that had been kept wet for ∼3 years were all laths and seemed to be exclusively Ia (see section 3.6 and [SI Figure S3](http://pubs.acs.org/doi/suppl/10.1021/acs.cgd.6b00981/suppl_file/cg6b00981_si_001.pdf)). A very small number of dried crystals were semiregular squat slabs, one of which gave diffraction showing two reciprocal lattices that had diffraction spots with similar relative intensities. We did not find any crystals that were single-component Ib. Reconstructed diffraction images showing slices of reciprocal space in the (0kl) plane for a lath-shaped crystal (UK data) and for a squat slab (ALS data) are shown in Figure 6. For the lath, reflections from the main lattice are clearly dominant, but there are additional weak Bragg peaks between some of the main reflections along  $c^*$ . For the slab, the "extra" reflections are considerably more intense. We initially suspected twinning, but these extra reflections could not be accounted for using any twin law. Temperature variation between 90 K and ∼250 K had little discernible effect on the relative intensities of the "extra" reflections. Attempts to interpret them as modulation-induced satellite peaks were futile. For data from solvent-free crystals that had such extra reflections, difference Fourier maps sometimes gave small but non-negligible (∼1−2 e·Å<sup>−</sup><sup>3</sup> ) peaks at a position well removed from the iodine atom in Ia. While this "ghost" peak could be interpreted as a minor disorder component for the iodine, such a disorder model does nothing to explain the presence of the extra reflections. In spite of the fact that all crystals gave sharp extinctions when rotated between crossed polarizers, the likelihood that these were hybrid crystals containing a variable amount of an unknown phase was compelling. At the ALS, this second phase indexed as Ib with a unit cell volume  $[651.3(2)$   $\AA$ <sup>3</sup>] that is significantly larger than that of Ia [615.95(5)  $\rm \AA^3$ ]. We suspected that Ib might be an epitaxially intergrown hydrate, similar to some uric acid crystals,[19,20](#page-9-0) but subsequent analysis proved it to be a true second polymorph.

3.6. Powder Diffraction from Solvent-Free I. Experimental powder diffraction patterns obtained from a bulk sample of pulverized crystals ([Figure 7\)](#page-6-0) gave a series of peaks that correspond very closely to peaks in a simulated powder pattern calculated from the crystal structure of Ia. A few small peaks, however, did not match up, and these anomalous peaks were reproducible across powders prepared from different batches of solvent-free crystals. A simulated powder pattern calculated from the crystal structure of Ib, however, readily explains these features. As with the single-crystal data, it is clear from these powder patterns that Ia is by far the dominant polymorph. A powder pattern from wet crystals that had sat undisturbed for ∼3 years showed virtually no Ib (see [SI Figure](http://pubs.acs.org/doi/suppl/10.1021/acs.cgd.6b00981/suppl_file/cg6b00981_si_001.pdf) [S4\)](http://pubs.acs.org/doi/suppl/10.1021/acs.cgd.6b00981/suppl_file/cg6b00981_si_001.pdf). There was no evidence of the mixed solvate form Ic in any experimental powder patterns.



Figure 6. Reciprocal-space-slice images reconstructed from a full set of diffraction frames for (a) lath-shaped and (b) slab-shaped crystals of solvent-free I, each showing the presence of two reciprocal lattices. In (a) reflections from one lattice are clearly dominant, whereas in (b) neither appears to be dominant.

3.7. Computational Modeling. A series of electronic structure calculations were performed to determine the influence of conformation on the different packing arrangements. DFT energies of the isolated molecules, using coordinates taken directly from the respective crystal structures, but with allowance for the hydrogen atoms to "relax" from their foreshortened X-ray crystal structure distances<sup>[39](#page-10-0)</sup> reveal the conformation in Ia to be less stable than that in Ib by ∼1.5 kcal/mol. On isolated molecules, the effect of torsional angle variation (i.e., about the N9−C1′ and C4′−C5′ bonds) was investigated through a series of geometry optimizations. The dihedral angle in question was fixed at the desired angle, while

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Figure 7. Qualitative comparison of powder diffraction patterns for solvent-free I. The black traces, (a) and (b), are experimental patterns from two different batches of crystals. The blue (c) and red (d) traces are simulated patterns calculated from the refined crystal structures of Ia and Ib respectively. The blue dotted vertical lines show the match up between experimental powder lines and those calculated from the crystal structure of Ia. The red vertical lines show the match up of two small features in the experimental powder patterns with peaks calculated from the crystal structure of Ib.

the rest of the molecule was allowed to relax. The angle was varied from  $-180^\circ$  to  $+180^\circ$  in  $5^\circ$  increments. Two optimizations were performed for each torsion, one based on the starting geometry of form Ia and the other starting from Ib.

3.7.1. Torsion about the Glycosidic Bond between Purine and Deoxyribose. For the deoxyribose−purine bond (i.e., torsion C4−N9−C1′−O1′), minimum energies were found for torsion angles of 75° and 175° when using the initial geometry of Ia (Figure 8). Starting from the molecular geometry of Ib, however, gives deeper potential wells for this torsion, with minima at −120° and 75°: the corresponding torsion angles in the crystal structures of Ia and Ib are  $-152.4(5)^{\circ}$  and −151(2)°, respectively ([Table 2\)](#page-3-0). The mismatch here between crystal structures and optimized isolated molecule conformations is due to the extensive intermolecular hydrogen-bonding network in the crystals. The implication, however, is that relaxation about the N9−C1′ bond leads to lower conformational energy if the iodine atom is placed as in Ib versus that in Ia.

3.7.2. Torsion of Iodine Relative to Deoxyribose. The torsion of the iodine relative to the ring oxygen of the deoxyribose group (i.e., I1−C5′−C4′−O1′) was investigated in a similar manner, yielding a minimum at 65° when starting from geometry Ia (Figure 9). Other local minima were found at −65° and 170°, both being around 1 kcal/mol higher in energy than the global minimum. The barrier for rotation for this bond was found to be around 6−8 kcal/mol. Starting from geometry Ib, the minimum at 65° is destabilized, while the minima at −65° and 175° become more stable, with the latter becoming the global minimum (Figure 9). There is good agreement between these optimized torsion angles and those measured in the crystal structures  $[61.2(5)^\circ$  and  $168.8(11)^\circ$  respectively],



Figure 8. Isolated-molecule conformational energies for torsion (C4− N9−C1′−O1′) around the purine-deoxyribose groups starting from the experimental geometries of Ia and Ib. The Ib conformation is able to relax to a lower-energy geometry than the Ia conformation.



Figure 9. Isolated-molecule conformational energies for torsion (I1− C5′−C4′−O1′) of the iodine relative to the deoxyribose group. The lowest energy conformer occurs at 175° when the starting model is Ib and is ∼1 kcal/mol lower than the lowest energy conformer (∼65°) when starting from Ia. For Ib, the trace is offset to lower energy by  $\sim$ 3 kcal/mol between about −130° and −160°: an effect caused by a change in the puckering of the deoxyribose ring. Computed geometries at each of the minima are shown in [SI Figure S5.](http://pubs.acs.org/doi/suppl/10.1021/acs.cgd.6b00981/suppl_file/cg6b00981_si_001.pdf)

while the two rotational degrees of freedom are shown to be dependent on one another. In the conformational analysis summarized in Figure 9, the torsions between about −130° and −160° for the Ib starting model are offset to lower energy by  $~\sim$ 3 kcal/mol. This unusual feature is due to an unexpected rearrangement of the deoxyribose ring conformation for the optimized structures. Interestingly, a similar rearrangement of the deoxyribose ring is evident in molecule  $Ic<sub>b</sub>$  of the mixed solvate ([Table 2](#page-3-0)), albeit at different angles.

3.7.3. Cohesive Energies between Layers of Molecules in Polymorphs **Ia** and **Ib**. To investigate the differences in <span id="page-7-0"></span>packing between Ia and Ib, specifically at the iodine−iodine layer interfaces, the interactions between molecules were calculated using a supermolecular approach. The energy of a pair of molecules (Figure 10) taken directly from the crystal



Figure 10. Pairs of molecules used to define iodine-involving interactions in (a) polymorph Ia and (b) polymorph Ib for pairwise interaction energy calculations. The structure fragments were taken directly from the refined crystal structures, and the hydrogen positions were allowed to relax.

structure was calculated after allowing hydrogen atoms to relax,<sup>[39](#page-10-0)</sup> using a dispersion-corrected DFT functional to properly describe the stabilizing dispersion intermolecular forces. In

addition to calculating the energy of the dimer, the energies of isolated individual molecules were calculated at the same computational level. The interaction energy,  $E_{int}$ , is given by the difference between the dimer and the sum of its isolated components. The interaction between layers in the c-direction involving iodine atoms was investigated through pairwise interactions, as well as the overall interaction between one molecule and its four nearest neighbors (Table 4). Both uncorrected and counterpoise corrected (to account for basis set superposition error) values are reported. The interdigitation in form Ia clearly leads to a more favorable interaction energy over form Ib due to closer intermolecular distances. The interaction energies obtained in the cluster approach do not completely match the sum of the pairwise interactions, where the electron density can be distributed to maximize the interaction between one pair, without decreasing an adjacent interaction, but are in fairly good agreement.

To describe the interactions in more detail, we used pairwise interaction energy decomposition analysis (PIEDA) $32$  to break down the overall interaction energy into separate electrostatic, exchange, charge transfer, and dispersion terms. The electrostatic energy is due to interactions of permanent partial charges on both molecules, while the charge transfer energy is due to the charge on one molecule inducing a temporary dipole on the second molecule. The exchange energy is the electron−electron repulsion between molecules, while the movement of electrons due to the above terms leads to temporary dipoles on both molecules, leading to the dispersion term. The PIEDA calculations were run on molecular pairs (vide supra, Figure 10) taken directly from the crystal structures, using the Møller− Plesset MP2 method and the same mixed basis sets as before. Once again, the interdigitating iodines give a stronger total interaction energy [\(Figure 11](#page-8-0)). In most of the pairs, the major stabilizing factor is dispersion, which is counteracted by the exchange repulsion. Electrostatic stabilization is also sizable. The exception is pair 1 of Ia, which shows a stronger electrostatic than dispersion energy. The minimal number of close contacts in Ib pair 2 results in overall weak interactions. Comparing these results with measurement of close intermolecular contacts in each pair (see [SI Table S2](http://pubs.acs.org/doi/suppl/10.1021/acs.cgd.6b00981/suppl_file/cg6b00981_si_001.pdf) for details), we can postulate that the increased electrostatic stabilization in Ia-1 may come from multiple close I···H interactions. Meanwhile, in addition to the stabilizing effects of halogen bonding $37$  and interhalogen interactions,<sup>[38](#page-10-0)</sup> interactions between electron rich atoms  $(I \cdots I$  and  $I \cdots O)$  contribute to stabilization by both dispersion and electrostatics observed in the PIEDA analysis of Ia-2.

## 4. DISCUSSION

4.1. Differences in Habit of Solvent-Free Crystals of (I). The vast majority of hybrid crystals were lath shaped, but a very small number were squat slabs. What could cause the observed differences in crystal shape? There are two obvious

Table 4. Pairwise and Cluster Interaction Energies between Closest Neighboring Molecules with Interactions Involving Iodine



 ${}^a$ Numbers in parentheses used counterpoise correction to account for the basis set superposition error.  ${}^b$ Specific molecule pairs are defined in Figure 10.

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Figure 11. MP2 Pairwise interaction energy decomposition analysis for close contacts between iodine-containing groups of adjacent molecules, after hydrogen relaxation. ES = electrostatic, EX = exchange, CT = charge transfer, DISP = dispersion, TOT = total interaction energy.

possibilities: either the squat slab was a fragment broken from a lath, or the presence of Ib impeded crystal growth, especially along b, thereby foreshortening the laths into slabs. The latter explanation is more attractive for several reasons. It is clear from [Figure 6](#page-5-0) that the laths contain considerably less Ib than the sole slab-shaped crystal that we analyzed in detail. On the basis of comparison of relative reflection intensities, a rough estimate for the fraction of Ib in typical lath-shaped crystals is 0−2%, while for the squat slab it was closer to 50%. Moreover, the squat slab did not look like a broken piece of lath. At the ALS, the squat slab was selected precisely because it looked so different from the rest of the crystals. Besides, if it were simply a fragment broken from a lath, it ought to have had a similar Ib content to the other laths that were analyzed in detail.

4.2. Compatibility of Unit Cells for Ia and Ib. The compatibility of hydrogen-bonding networks ([Figure 2](#page-3-0)) and the incompatibility of iodine-containing layers [\(Figure 3\)](#page-4-0) suggest that the interface between conjoined polymorphs consists of a hybrid double layer in the ab plane, composed of a layer of Ia hydrogen bonding with a layer of Ib. Epitaxy solely in the ab plane agrees with the evidence from diffraction: all "extra" reflections (e.g., [Figure 6](#page-5-0)) occur along directions parallel to  $c^*$ , which is perpendicular to the  $ab$  plane. The  $a$ -axis and  $b$ -axis lengths of Ia are ∼1% and ∼0.2% longer than those of Ib, leading to a mismatch of ∼1% for the areas of the ab faces of Ia and Ib. While these differences are small, the cumulative effect over several repeat distances rapidly lead to a substantial discrepancy (Figure 12), which could plausibly poison the growing crystal surface, thereby impeding crystal growth in the a and/or b-axis directions. This is consistent with the laths having a relatively low proportion of Ib and the squat slab having comparatively higher Ib content. The line of reasoning here is inspired by the crystal habit modification work of Leiserowitz and  $\text{Lahav}^{40,41}$  $\text{Lahav}^{40,41}$  $\text{Lahav}^{40,41}$  in which "tailored" dopant molecules were shown to inhibit crystal growth in specific directions.

4.3. Conformational Energies and Implications for Crystal Growth. For isolated molecules, the computed overall conformation energy for Ia is higher than that of Ib. The iodine atom torsions with the deoxyribose ring can be defined relative to O1′ and to C3′. In Ia, the former is gauche and the latter is anti, while in Ib these assignments are reversed. In Ia, the angles [\(Table 2](#page-3-0)) are closer to ideal values than Ib, but "ideal values" make no account of the crystal-packing environment in the solid state. For isolated molecules, although conformational energies for the two torsional degrees of freedom are coupled, computed energies suggest a ∼1.5 kcal/mol advantage for the Ib conformer. This of course ignores neighboring molecules, and so need not be relevant in the crystal or even in solution. In the crystalline state, Ia is clearly dominant, which combined with its higher density (2.120 vs 2.005 Mg m<sup>-3</sup> for Ia and Ib respectively) suggests it is thermodynamically more stable at room temperature. This makes intuitive sense given the striking difference in how the iodine layers interact in Ia vs Ib. Pairwise and cluster interaction energies for model layer fragments strongly favor the interdigitated Ia over the abutted Ib. What then to make of Ib? Since the ab layers are extensively hydrogen bonded in two directions it seems likely that fragments of such layers exist in solution. Such layers would be expected to look more like layers of Ib than Ia because the **Ib** form layers are ∼1.1% denser (*ab* = 37.66 vs 38.11 Å<sup>2</sup>) and the molecules have a more favorable conformation. When such a fragment joins the growing crystal there would be a driving force for the  $-CH_2I$  groups to "zip up" to give form Ia, but sometimes that process could fail, leading to domains of form Ib. In this scenario, Ib would be a metastable kinetic product. $42$ Efficient conversion of Ib to Ia could thus be dependent on the presence of solvent, and at least for wet crystals, there was a tendency for the presence of Ib to diminish over time.



Figure 12. An overlay of 11 repeats along the a-axis for Ia and Ib, projected perpendicular to ac, based on a least-squares fit of atoms in the purine groups in the central pairs of overlaid molecules. The unit cells of Ia and Ib tilt to the left and right, respectively. The mismatch along the a-direction rapidly accumulates and could plausibly poison the surface, impeding growth in the  $a$  and/or  $b$ -directions, thereby foreshortening lath-shaped crystals into squat slabs for high Ib content crystals.

# <span id="page-9-0"></span>5. SUMMARY

Solvent-free crystals of 5′-deoxy-5′-iodo guanosine coexist as hybrids of two conformational polymorphs that differ primarily in the orientation of the 5′-iodo substituent relative to the deoxyribose ring. The guanosine groups in each polymorph are in similar enough orientations to allow epitaxial intergrowth of the two forms. Although the more common polymorph could be obtained in essentially pure form, the less common polymorph was only found in hybrid crystals containing both polymorphs. Computational analysis suggests that for isolated molecules, the less common polymorph geometry is energetically more favorable. In the crystalline state, however, this leads to a much less favorable packing arrangement between layers of exposed iodine substituents. The implication is that the less common polymorph is a kinetic product, while the more common polymorph is the thermodynamically more stable crystalline form at room temperature.

# ■ ASSOCIATED CONTENT

#### **6** Supporting Information

The Supporting Information is available free of charge on the [ACS Publications website](http://pubs.acs.org) at DOI: [10.1021/acs.cgd.6b00981.](http://pubs.acs.org/doi/abs/10.1021/acs.cgd.6b00981)

Figure S1. A photograph of typical lath-shaped crystals of I. Figure S2. A packing plot for the mixed solvate Ic viewed down the  $(a)$  a axis and  $(b)$  b axis. Figure S3. Reciprocal-space-slice images for the (0kl) and (h0l) layers for a solvent-free crystal that had aged in mother liquor for ∼3 years. Figure S4. Qualitative powder diffraction pattern for pulverized solvent-free crystal that had aged in mother liquor for ∼3 years. Figure S5. Computed geometries at the minimum points shown in [Figure 9](#page-6-0). Table S1. Hydrogen bond parameters for the mixed water/methanol solvate Ic. Table S2. Closecontact distances (Å) for different types of atom−atom interactions between pairs of molecules defined in [Figure](#page-7-0) [10](#page-7-0) of the main manuscript ([PDF](http://pubs.acs.org/doi/suppl/10.1021/acs.cgd.6b00981/suppl_file/cg6b00981_si_001.pdf))

# Accession Codes

CCDC [1486096](https://summary.ccdc.cam.ac.uk/structure-summary?pid=ccdc:1486096&id=doi:10.1021/acs.cgd.6b00981)−[1486098](https://summary.ccdc.cam.ac.uk/structure-summary?pid=ccdc:1486098&id=doi:10.1021/acs.cgd.6b00981) contain the supplementary crystallographic data for this paper. These data can be obtained free of charge via [www.ccdc.cam.ac.uk/data\\_request/cif](http://www.ccdc.cam.ac.uk/data_request/cif), or by emailing data request@ccdc.cam.ac.uk, or by contacting The Cambridge Crystallographic Data Centre, 12 Union Road, Cambridge CB2 1EZ, UK; fax: +44 1223 336033.

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## Author Contributions

S.R.P. and K.J.G. performed crystallographic analyses, K.J.T. conducted quantum-chemical calculations, and E.J.B. was responsible for synthesis and crystal growth. The manuscript includes contributions from all authors, and all authors have given approval to the final version of the manuscript.

# Notes

The authors declare no competing financial interest.

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